

# A Macroscopic Description of Gauging — as a special Morita equivalence

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# Outline

Gauging is usually done microscopically in field theories or lattice models to give a dual (gauged) theory.

We give a macroscopic/categorical understanding.

The following important data are needed for gauging:

- The symmetry assignment of the theory to be gauged  
**a monoidal functor**
- The precise way to perform gauging: which symmetry defects are to be condensed and how  
**a module category** or a condensation algebra of symmetry defects

The remaining things are just computing centers!

# Outline

- 1 Symmetry assignment
- 2 Center functor, Zipping, Unzipping and “sandwich”
- 3 Abstract gauging, gauging prescription and reversibility
- 4 Examples

# Symmetry assignment

The traditional way of assigning a symmetry  $G$  to a quantum system  $V$  is by specifying a group representation:

$$\rho : G \rightarrow GL(V) \subset \text{End}(V).$$

- $G$  is an abstract (indexing) group.
- $\text{End}(V)$  is all linear operators on  $V$ .
- $\rho$  is a group homomorphism a map that preserves multiplication.

It is important to separate the abstract multiplication rule from concrete operators. Even if  $G$  and  $V$  are the same, different  $\rho$  should be considered as different physical symmetries.

To proceed, we need to abandon the idea that

~~$G$  acts externally on  $V$ .~~

# Symmetry assignment

Generalized symmetry are formed by topological operators.

- If we collect all the possible topological operators in a physical system in  $n+1$ D, the corresponding mathematical description is a fusion  $n$ -category  $\mathcal{A}$ .
- The abstract (indexing) generalized symmetry is also a fusion  $n$ -category  $\mathcal{T}$ .

## Definition

A symmetry assignment is a higher linear monoidal functor

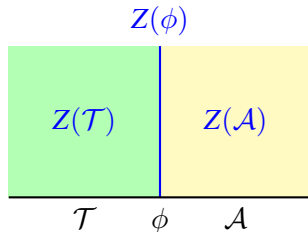
$$\phi : \mathcal{T} \rightarrow \mathcal{A}.$$

# Center functor

The center functor plays a central role in gauging.

A useful physical picture is to think center computes bulk.

- Center of a fusion  $n$ -category  $\mathcal{C}$ :  $Z(\mathcal{C}) = \text{Fun}_{\mathcal{C}|\mathcal{C}}(\mathcal{C}, \mathcal{C})$ .
- Center of  $\mathcal{C}$ - $\mathcal{D}$ -bimodule  $\mathcal{M}$ :  $Z_{\mathcal{C}|\mathcal{D}}(\mathcal{M}) = \text{Fun}_{\mathcal{C}|\mathcal{D}}(\mathcal{M}, \mathcal{M})$ .
- Center of  $\phi : \mathcal{T} \rightarrow \mathcal{A}$ :  $Z(\phi) := Z_{\mathcal{T}|\mathcal{A}}(\phi\mathcal{A})$ , where  $\phi\mathcal{A}$  is the  $\mathcal{T}$ - $\mathcal{A}$ -bimodule  $\mathcal{A}$  with left action given by  $\phi$ .



Center depicted as bulk.

# Physical meaning of $Z(\phi)$

Informally speaking,  $Z(\phi)$  is the centralizer or commutant of  $\phi(\mathcal{T})$  in  $\mathcal{A}$ :

$$Z(\phi) \sim \{X \in \mathcal{A} \mid X \otimes \phi(s) \cong \phi(s) \otimes X, \forall s \in \mathcal{T}\}.$$

$Z(\phi)$  is the topological operators that are invariant under the action of  $\phi(\mathcal{T})$ , recall that a unitary symmetry acts on operators by conjugation  $U_g X U_g^\dagger$ , which may be called the charge/representation category.

Example: for the forgetful functor  $\phi : n\mathbf{Vec}_G \rightarrow n\mathbf{Vec}$ ,  
 $Z(\phi) = n\mathbf{Rep}G$ .





# Unzipping

The other direction, we call Unzipping:

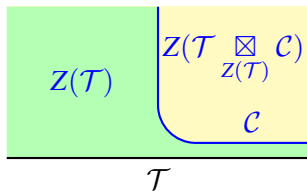
Sandwich  $\mathcal{T} \boxtimes_{Z(\mathcal{T})} \mathcal{C} \cong \mathcal{A}$  implies a natural symmetry assignment:

$$\eta_{\mathcal{T}, \mathcal{C}} : \mathcal{T} \rightarrow \mathcal{T} \boxtimes_{Z(\mathcal{T})} \mathcal{C}$$

$$s \mapsto s \boxtimes_{Z(\mathcal{T})} \mathbf{1}_{\mathcal{C}}$$

The center of this symmetry assignment is exactly  $\mathcal{C}$ ,

$$\mathcal{C} \cong Z(\eta_{\mathcal{T}, \mathcal{C}}).$$



# Abstract gauging

The ways of gauging should be first specified for the abstract symmetry  $\mathcal{T}$ :

## Definition

An abstract gauging of the abstract symmetry  $\mathcal{T}$  is an indecomposable right  $\mathcal{T}$ -module  $\mathcal{K}$ .

Equivalently, an abstract gauging can be specified by an algebra  $K$  in  $\mathcal{T}$  such that  $\text{LMod}_K(\mathcal{T}) \cong \mathcal{K}$ . The symmetry defects contained in the algebra  $K$  are to be “summed over” or “condensed”.

# Abstract gauging

## Definition

An abstract gauging of the abstract symmetry  $\mathcal{T}$  is an indecomposable right  $\mathcal{T}$ -module  $\mathcal{K}$ .

Denote by  $\mathcal{T}_{\mathcal{K}}^{\vee} := Z_{n\text{Vec}|\mathcal{T}}(\mathcal{K}) = \text{Fun}_{n\text{Vec}|\mathcal{T}}(\mathcal{K}, \mathcal{K})$  the center of  $\mathcal{K}$ , called the gauge symmetry or dual symmetry (of  $\mathcal{T}$  with respect to  $\mathcal{K}$ ).

$\mathcal{T}_{\mathcal{K}}^{\vee}$  is Morita equivalent to  $\mathcal{T}$  by definition.

# Abstract gauging

## Definition

An abstract gauging of the abstract symmetry  $\mathcal{T}$  is an indecomposable right  $\mathcal{T}$ -module  $\mathcal{K}$ .

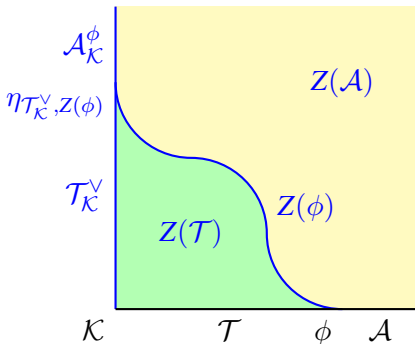
If  $n\mathbf{Vec}$  is a  $\mathcal{T}$ -module,  $\mathcal{T}$  is a local fusion  $n$ -category (free of 't Hooft anomaly) and choosing  $\mathcal{K} = n\mathbf{Vec}$  corresponds to the ordinary gauging, i.e., “summing over” all possible symmetry defects.

In general,  $\mathcal{K}$  can be larger and corresponds to a partial gauging where only a subset of symmetry defects, instead of all, are “summed over”.

# Gauging prescription

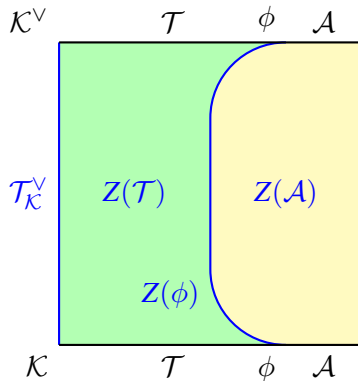
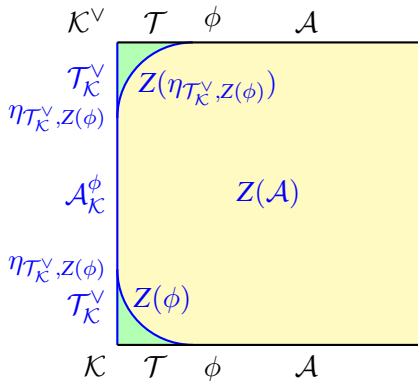
## Definition

Given a theory  $\mathcal{A}$ , we need to specify its symmetry  $\phi : \mathcal{T} \rightarrow \mathcal{A}$ , and an abstract gauging  $\mathcal{K}$  of  $\mathcal{T}$ , and then the gauged theory is  $\mathcal{A}_{\mathcal{K}}^{\phi} := \mathcal{T}_{\mathcal{K}}^{\vee} \boxtimes_{Z(\mathcal{T})} Z(\phi) \cong Z_{n\mathbf{Vec}|\mathcal{A}}(\mathcal{K} \boxtimes_{\mathcal{T}} \phi \mathcal{A})$ .

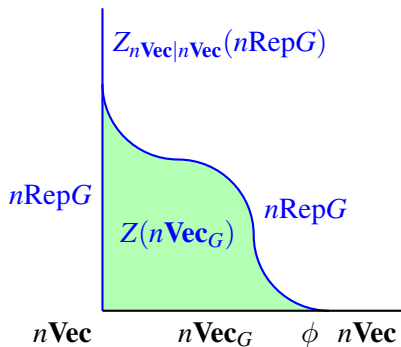


By Unzipping,  $\mathcal{T}_{\mathcal{K}}^{\vee}$  is a natural symmetry of the gauged theory  $\mathcal{A}_{\mathcal{K}}^{\phi}$ , with the same charge category  $Z(\phi)$  as the original theory.

# Gauging is reversible (Morita equivalence)

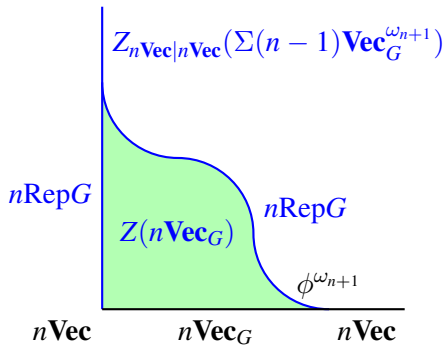


# Ordinary gauge theory



$\phi : n\text{Vec}_G \rightarrow n\text{Vec}$  is the forgetful functor. The gauged theory is the ordinary  $G$ -gauge theory.

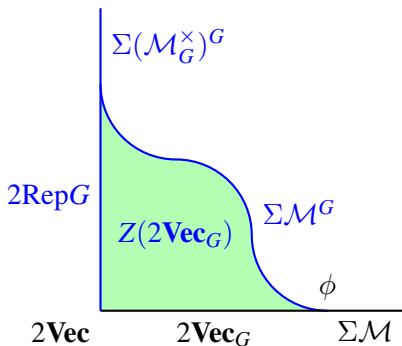
# Dijkgraaf-Witten gauge theory



$\phi^{\omega_{n+1}} : n\mathbf{Vec}_G \rightarrow n\mathbf{Vec}$  is twisted by an  $(n+1)$ -cocycle  $\omega_{n+1}$ , describing a SPT phase. The gauged theory is the Dijkgraaf-Witten  $G$ -gauge theory.



# Gauging 2+1D SET phases



Take a UMTC  $\mathcal{M}$  with a  $G$ -action. It can be gauged iff there exists a  $G$ -crossed braided extension  $\mathcal{M}_G^\times$ , and in the 2-categorical language, iff there exists a monoidal functor  $\phi : 2\mathbf{Vec}_G \rightarrow \Sigma\mathcal{M}$ .

ENO arXiv:0909.3140, BBCW arXiv:1410.4540

$\mathcal{M}^G$ , the equivariantization of  $\mathcal{M}$ , is a UMTC over  $\mathbf{Rep}G$ ;  $(\mathcal{M}_G^\times)^G$  is a minimal modular extension of  $\mathcal{M}^G$ .

LKW arXiv:1602.05936,1602.05946

## Other Examples

- Partial gauging: take a subgroup  $L \subset G$ . The abstract gauging  $\mathcal{K} = \text{LMod}_{(n-1)\text{Vec}_L}(n\text{Vec}_G)$  is equivalent to partially gauging the subgroup  $L$ .
- Higher gauging: we allow  $Z(\mathcal{A})$  to be nontrivial, in which case gauging  $\mathcal{A}$  is a higher gauging inside a nontrivial bulk.
- “Lower” gauging: we can as well gauge the bulk theory  $\Sigma Z(\mathcal{A})$ . A typical example is gauging fermion number parity.

## Conclusion and Outlook

- Symmetry assignment  $\phi : \mathcal{T} \rightarrow \mathcal{A}$  is equivalent to sandwich  $\mathcal{T} \boxtimes_{Z(\mathcal{T})} Z(\phi) \cong \mathcal{A}$ .
- Abstract gauging  $\mathcal{K}$  is a right  $\mathcal{T}$ -module.
- Gauged theory is the center of right  $\mathcal{A}$ -module  $\mathcal{K} \boxtimes_{\mathcal{T}} \phi \mathcal{A}$  and thus Morita equivalent to  $\mathcal{A}$ .
- Our framework covers all known variants of gauging, and may be used to discover unknown ones

Thanks for attention!